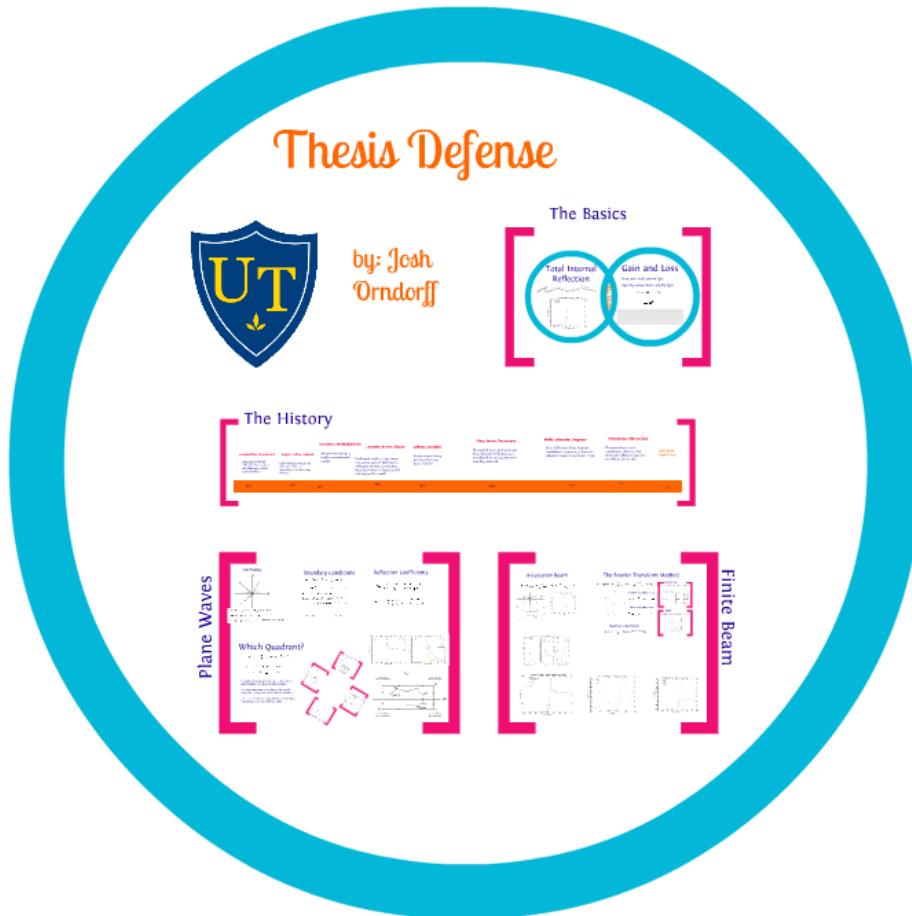


Amplified Total Internal Reflection at the Surface of a Gain Medium



Special thanks to Dr. Deck,
Dr. Karpov, and Dr. Bagley

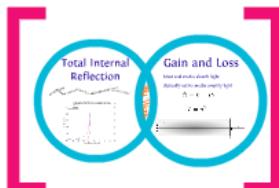
Amplified Total Internal Reflection at the Surface of a Gain Medium

Thesis Defense



by: Josh
Orndorff

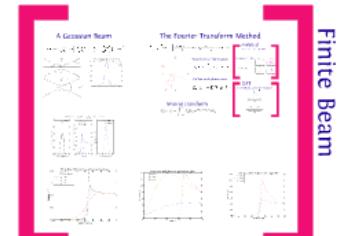
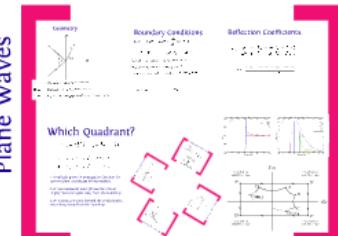
The Basics



The History

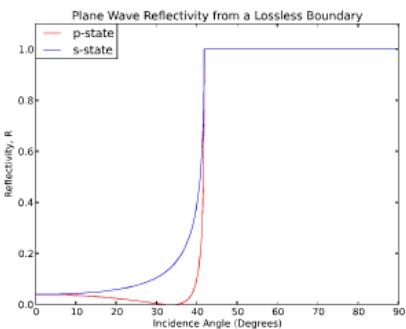
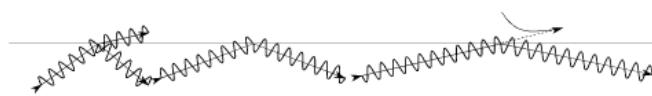


plane Waves



Finite Beam

Total Internal Reflection



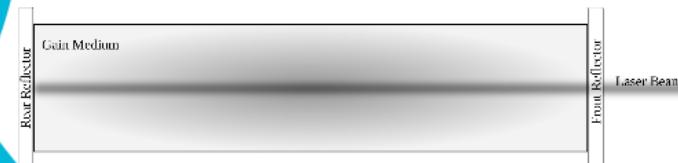
Gain and Loss

Most real media absorb light

Optically active media amplify light

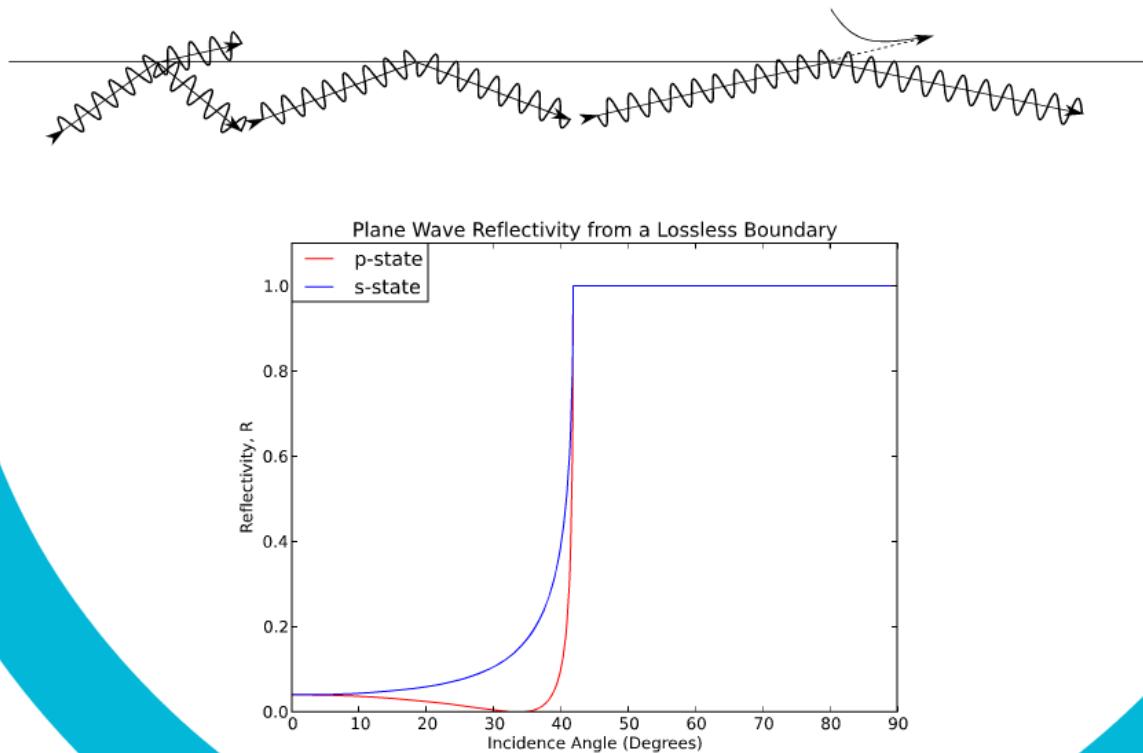
$$\tilde{n} = n - i\gamma$$

$$\tilde{\epsilon} = \tilde{n}^2$$



Why Work

Total Internal Reflection



WOW!



Ga

Most re

Optically

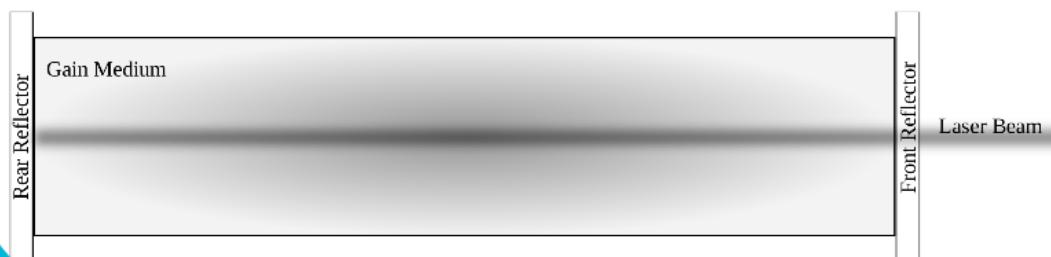
Gain and Loss

Most real media absorb light

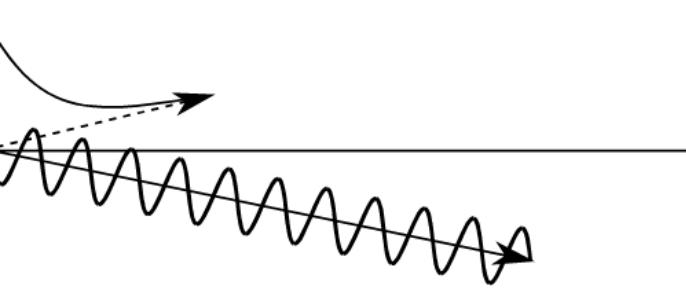
Optically active media amplify light

$$\tilde{n} = n - i\gamma$$

$$\tilde{\epsilon} = \tilde{n}^2$$



on



Most real m

Optically act

\tilde{n}

Gain Medium

ar Reflector



The History

Koester Fiber Experiment	Romanov, Shakhidzhanov	Lebedev, Volkov, Kogan	Callary, Carniglia	Plotz, Simon, Tucciarone	Willis, Schneider, Hagness	Mansuripur, Mansuripur
Detected Amplified reflection from a gain-clad fiber optic cable experimentally	Wrote first theory to explain experimental results	Performed another experiment measuring gain of 1000 from a reflecting surface, concluding that the Romanov theory could not explain the result	Wrote a new theory treating the three-layer problem	Theoretical work demonstrated that enhanced reflection was possible when surface plasmon coupling occurred	Finite Difference Time Domain simulations supporting enhanced reflection beyond the critical angle	Theoretical work and simulations claiming that enhanced reflection was not possible at any angles

1966

1972

1972

1973

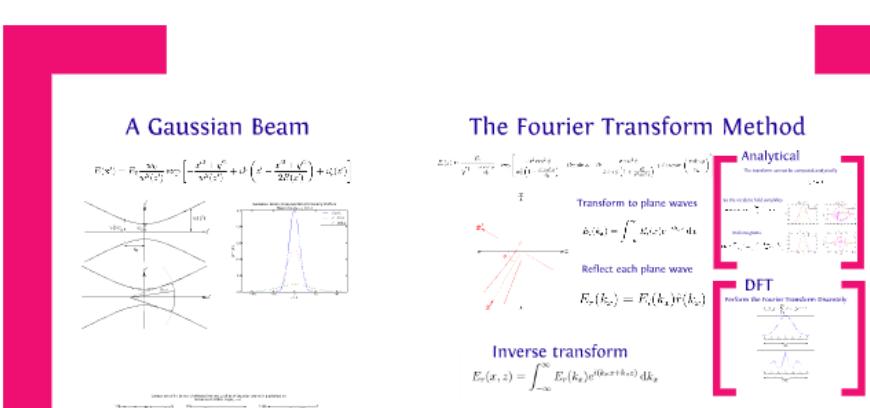
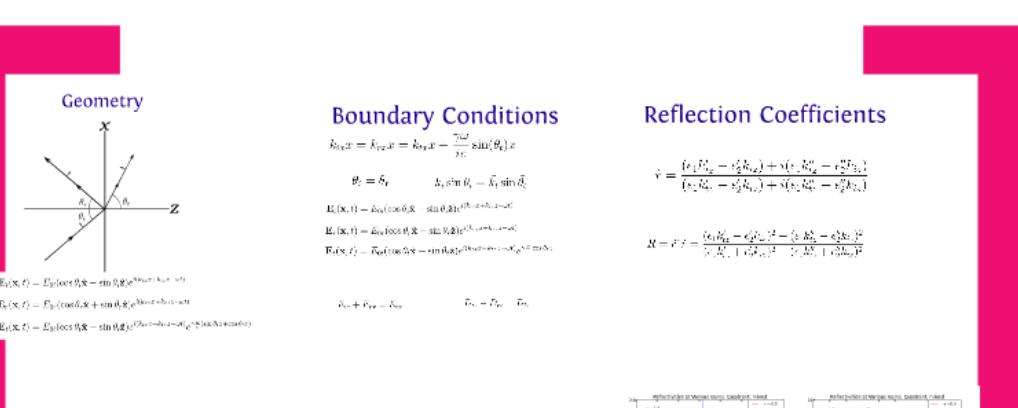
1976

1979

2003

2012

2013



Koester Fiber Experiment

Detected Amplified
reflection from a gain-
clad fiber optic cable
experimentally

1966

ment

Kogan, Volkov, Lebedev

Experimentally measured
gain of 25 from an
optically active reflecting
surface

1972

1972

ov, Lebedev

measured
an
reflecting

Romanov, Shakhidzhanov

Le

Wrote first theory to
explain experimental
results

Performed
measuring
reflecting
that the R
not explain

1972

zhanov

Lebedev, Volkov, Kogan

Performed another experiment measuring gain of 1000 from a reflecting surface, concluding that the Romanov theory could not explain the result

1973

Kogan

Callary, Carniglia

Wrote a new theory
treating the three-
layer problem

1976

Plotz, Simon, Tucciarone

Theoretical work demonstrated that enhanced reflection was possible when surface plasmon coupling occurred

Willis, Schneider, Hagness

Finite Difference Time Domain
simulations supporting enhanced
reflection beyond the critical angle

2003

Mansuripur, Mansuripur

Theoretical work and simulations claiming that enhanced reflection was not possible at any angles

Josh's
Degree

2012

20

, Mansuripur

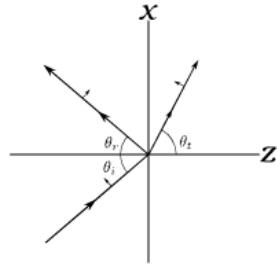
k and
ming that
tion was not
angles

Josh's Master
Degree Thesis

2013

Plane Waves

Geometry



$$\mathbf{E}_i(\mathbf{x}, t) = E_{0i}(\cos \theta_i \hat{\mathbf{x}} - \sin \theta_i \hat{\mathbf{z}}) e^{i(k_{ix}x + k_{iz}z - \omega t)}$$

$$\mathbf{E}_r(\mathbf{x}, t) = E_{0r}(\cos \theta_r \hat{\mathbf{x}} + \sin \theta_r \hat{\mathbf{z}}) e^{i(k_{rx}x + k_{rz}z - \omega t)}$$

$$\mathbf{E}_t(\mathbf{x}, t) = E_{0t}(\cos \theta_t \hat{\mathbf{x}} - \sin \theta_t \hat{\mathbf{z}}) e^{i(k_{tx}x + k_{tz}z - \omega t)} e^{\gamma_c^w (\sin \theta_t x + \cos \theta_t z)}$$

Boundary Conditions

$$k_{ix}x = k_{rx}x = k_{tx}x + \frac{\gamma\omega}{ic} \sin(\theta_t)x$$

$$\theta_i = \theta_r \quad k_i \sin \theta_i = k_t \sin \theta_t$$

$$\mathbf{E}_i(\mathbf{x}, t) = E_{0i}(\cos \theta_i \hat{\mathbf{x}} - \sin \theta_i \hat{\mathbf{z}}) e^{i(k_{ix}x + k_{iz}z - \omega t)}$$

$$\mathbf{E}_r(\mathbf{x}, t) = E_{0r}(\cos \theta_r \hat{\mathbf{x}} + \sin \theta_r \hat{\mathbf{z}}) e^{i(k_{rx}x + k_{rz}z - \omega t)}$$

$$\mathbf{E}_t(\mathbf{x}, t) = E_{0t}(\cos \theta_t \hat{\mathbf{x}} - \sin \theta_t \hat{\mathbf{z}}) e^{i(k_{tx}x + k_{tz}z - \omega t)} e^{\gamma_c^w \cos \theta_t z}$$

$$E_{ix} + E_{rx} = E_{tx}$$

$$D_{iz} + D_{rz} = D_{tz}$$

Reflection Coefficients

$$\tilde{r} = \frac{(\epsilon_1 k'_{tz} - \epsilon'_2 k_{iz}) + i(\epsilon_1 k''_{tz} - \epsilon''_2 k_{iz})}{(\epsilon_1 k'_{iz} + \epsilon'_2 k_{iz}) + i(\epsilon_1 k''_{iz} + \epsilon''_2 k_{iz})}$$

$$R = \tilde{r}^* \tilde{r} = \frac{(\epsilon_1 k'_{tz} - \epsilon'_2 k_{iz})^2 + (\epsilon_1 k''_{tz} - \epsilon''_2 k_{iz})^2}{(\epsilon_1 k'_{iz} + \epsilon'_2 k_{iz})^2 + (\epsilon_1 k''_{iz} + \epsilon''_2 k_{iz})^2}$$

Which Quadrant?

$$\tilde{k}_{tz}^2 = k_0^2 (\epsilon'_2 - \epsilon_1 \sin^2 \theta_i + i\epsilon''_2)$$

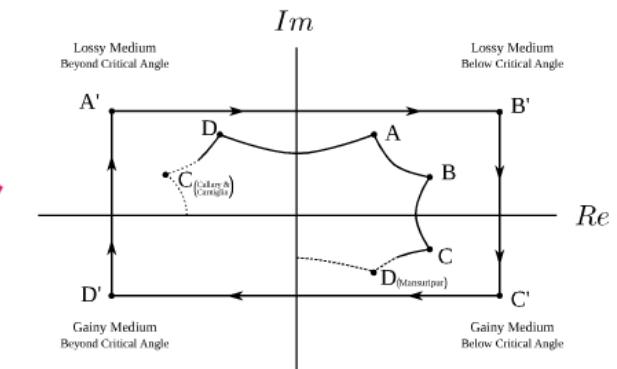
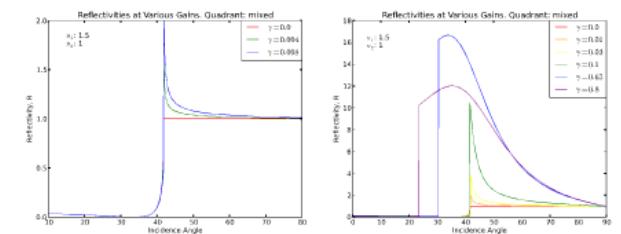
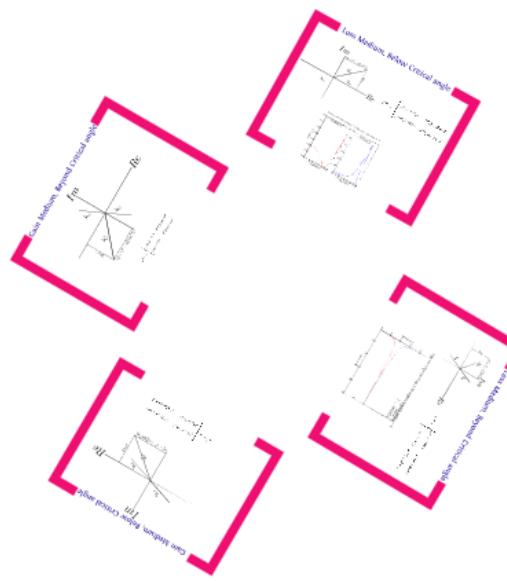
$$k'_{iz} = 1/\sqrt{2} \sqrt{\epsilon'_2 - \epsilon_1 \sin^2 \theta_i - \sqrt{(\epsilon'_2 - \epsilon_1 \sin^2 \theta_i)^2 - \epsilon''_2^2}}$$

$$k''_{iz} = \pm 1/\sqrt{2} \sqrt{\epsilon_1 \sin^2 \theta_i - \epsilon'_2 - \sqrt{(\epsilon'_2 - \epsilon_1 \sin^2 \theta_i)^2 - \epsilon''_2^2}}$$

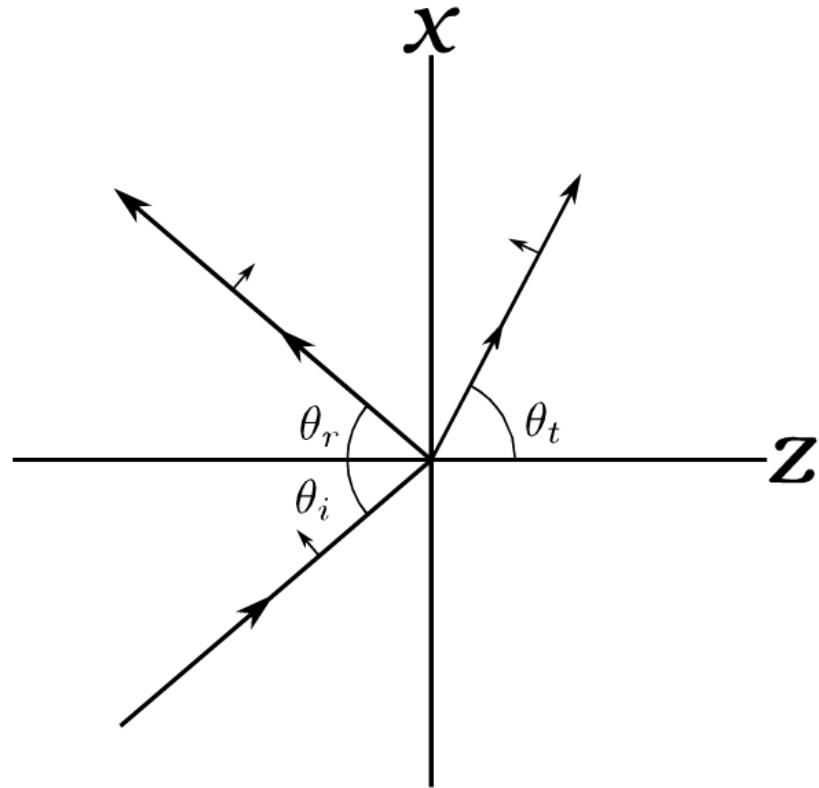
1. Amplitude grows in propagation direction for gain medium, and decays in loss medium.

2. A non-evanescent wave (below the critical angle) must propagate away from the boundary.

3. An evanescent wave (beyond the critical angle) must decay away from the boundary.



Geometry



$$\mathbf{E}_i(\mathbf{x}, t) = E_{0i}(\cos \theta_i \hat{\mathbf{x}} - \sin \theta_i \hat{\mathbf{z}}) e^{i(k_{ix}x + k_{iz}z - \omega t)}$$

$$\mathbf{E}_r(\mathbf{x}, t) = E_{0r}(\cos \theta_r \hat{\mathbf{x}} + \sin \theta_r \hat{\mathbf{z}}) e^{i(k_{rx}x + k_{rz}z - \omega t)}$$

$$\mathbf{E}_t(\mathbf{x}, t) = E_{0t}(\cos \theta_t \hat{\mathbf{x}} - \sin \theta_t \hat{\mathbf{z}}) e^{i(k_{tx}x + k_{tz}z - \omega t)} e^{\gamma \frac{\omega}{c} (\sin \theta_t x + \cos \theta_t z)}$$

Boundary Conditions

$$k_{ix}x = k_{rx}x = k_{tx}x + \frac{\gamma\omega}{ic} \sin(\theta_t)x$$

$$\theta_i = \theta_r \quad k_i \sin \theta_i = \tilde{k}_t \sin \tilde{\theta}_t$$

$$\mathbf{E}_i(\mathbf{x}, t) = E_{0i} (\cos \theta_i \hat{\mathbf{x}} - \sin \theta_i \hat{\mathbf{z}}) e^{i(k_{ix}x + k_{iz}z - \omega t)}$$

$$\mathbf{E}_r(\mathbf{x}, t) = E_{0r} (\cos \theta_r \hat{\mathbf{x}} + \sin \theta_r \hat{\mathbf{z}}) e^{i(k_{rx}x + k_{rz}z - \omega t)}$$

$$\mathbf{E}_t(\mathbf{x}, t) = E_{0t} (\cos \theta_t \hat{\mathbf{x}} - \sin \theta_t \hat{\mathbf{z}}) e^{i(k_{tx}x + k_{tz}z - \omega t)} e^{\gamma \frac{\omega}{c} \cos \theta_t z}$$

$$E_{ix} + E_{rx} = E_{tx}$$

$$D_{iz} + D_{rz} = D_{tz}$$

$$\cos \theta_t z)$$

Reflection Coefficients

$$\tilde{r} = \frac{(\epsilon_1 k'_{tz} - \epsilon'_2 k_{iz}) + i(\epsilon_1 k''_{tz} - \epsilon''_2 k_{iz})}{(\epsilon_1 k'_{tz} + \epsilon'_2 k_{iz}) + i(\epsilon_1 k''_{tz} + \epsilon''_2 k_{iz})}$$

$$R = \tilde{r}^* \tilde{r} = \frac{(\epsilon_1 k'_{tz} - \epsilon'_2 k_{iz})^2 + (\epsilon_1 k''_{tz} - \epsilon''_2 k_{iz})^2}{(\epsilon_1 k'_{tz} + \epsilon'_2 k_{iz})^2 + (\epsilon_1 k''_{tz} + \epsilon''_2 k_{iz})^2}$$

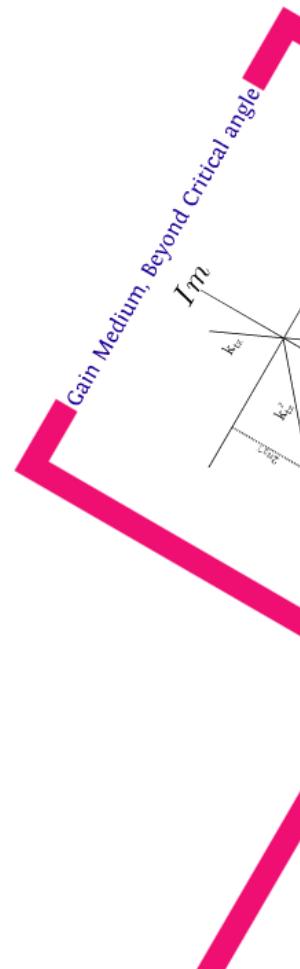
Which Quadrant?

$$\tilde{k^2}_{tz} = k_0^2 \left(\epsilon'_2 - \epsilon_1 \sin^2 \theta_i + i\epsilon''_2 \right)$$

$$k'_{tz} = \pm \frac{k_0}{\sqrt{2}} \sqrt{\epsilon'_2 - \epsilon_1 \sin^2 \theta_i + \sqrt{(\epsilon'_2 - \epsilon_1 \sin^2 \theta_i)^2 + \epsilon''_2^2}}$$

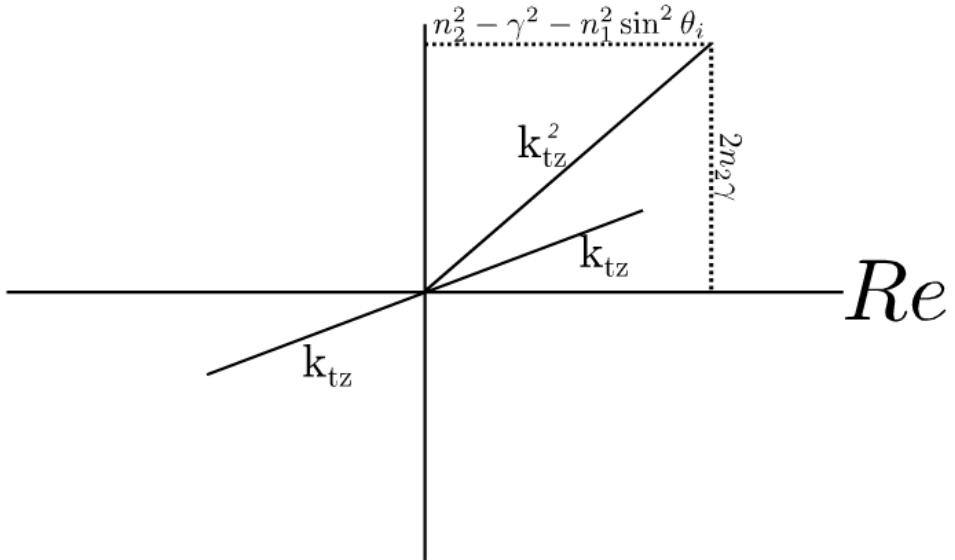
$$k''_{tz} = \pm \frac{k_0}{\sqrt{2}} \sqrt{\epsilon_1 \sin^2 \theta_i - \epsilon'_2 + \sqrt{(\epsilon'_2 - \epsilon_1 \sin^2 \theta_i)^2 + \epsilon''_2^2}}$$

1. Amplitude grows in propagation direction for gain medium, and decays in loss medium.
2. A non-evanescent wave (below the critical angle) must propagate away from the boundary.
3. An evanescent wave (beyond the critical angle) must decay away from the boundary.



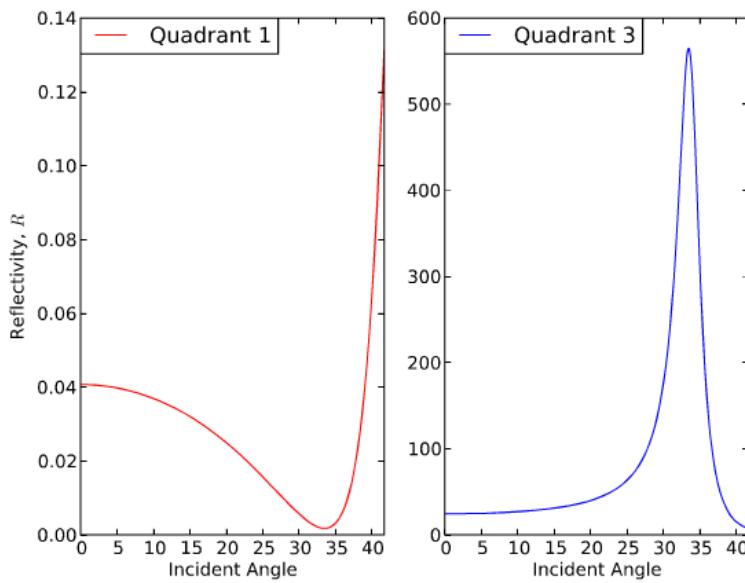
Loss Medium, Below Critical angle

Im

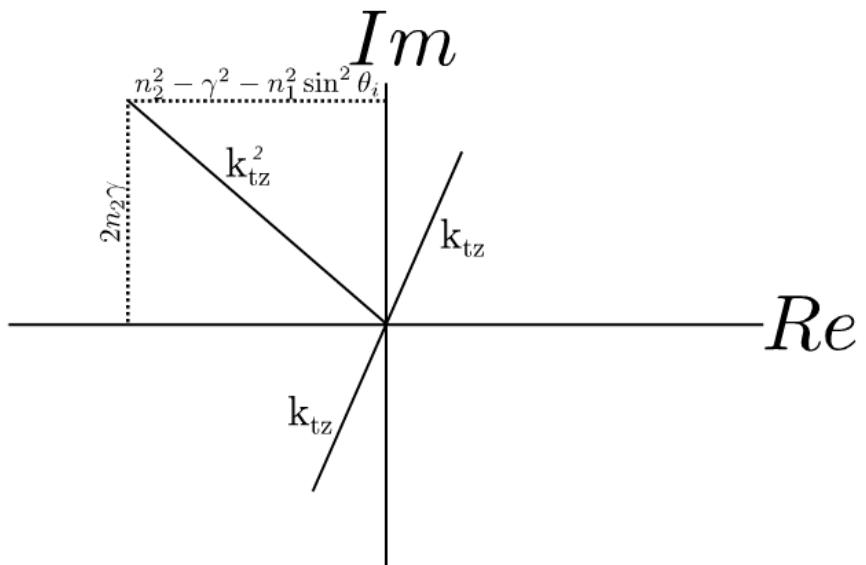


$$e^{i\tilde{k}_{tz}} = \begin{cases} e^{i|k'_{tz}|z - |k''_{tz}|z} & 1^{st} Quadrant \\ e^{-i|k'_{tz}|z + |k''_{tz}|z} & 3^{rd} Quadrant \end{cases}$$

Plane wave reflectivity options, R , for incidence on a lossy medium, $\theta_i < \theta_c$

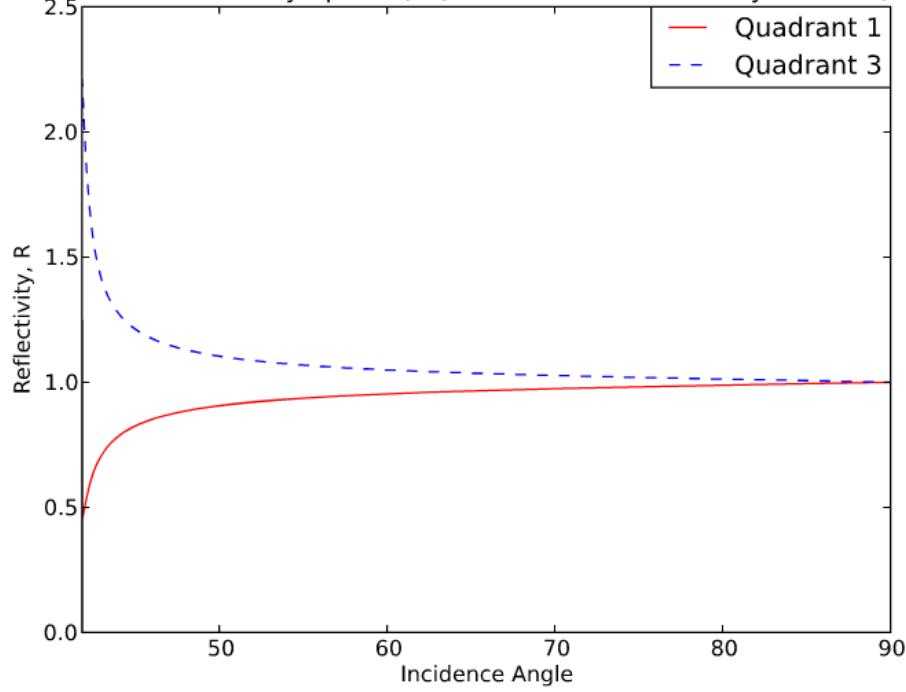


Loss Medium, Beyond Critical angle



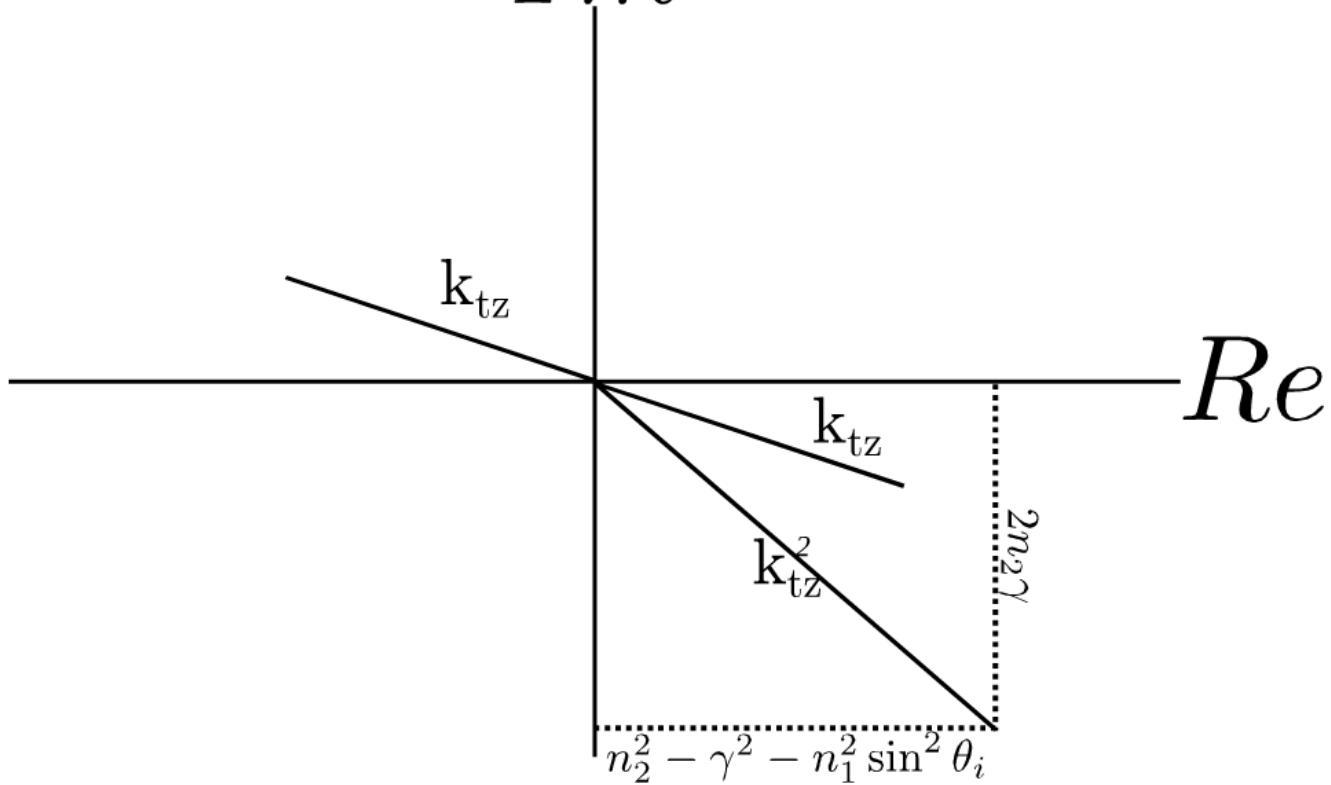
$$e^{i\tilde{k}_{tz}} = \begin{cases} e^{i|k'_{tz}|z - |k''_{tz}|z} & 1^{st} Quadrant \\ e^{-i|k'_{tz}|z + |k''_{tz}|z} & 3^{rd} Quadrant \end{cases}$$

Plane wave reflectivity options, R , for incidence on a lossy medium, $\theta_i < \theta_c$



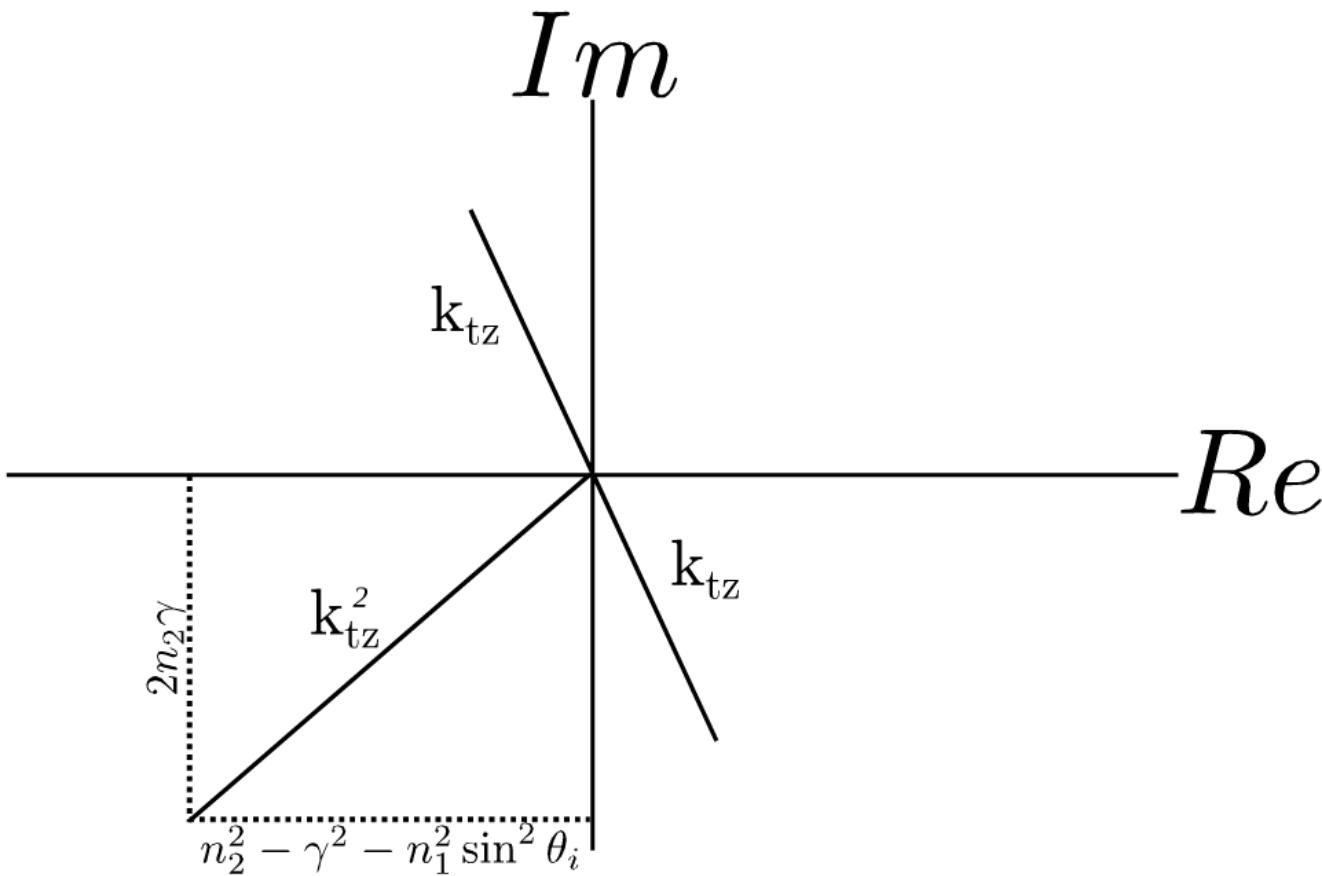
Gain Medium, Below Critical angle

Im



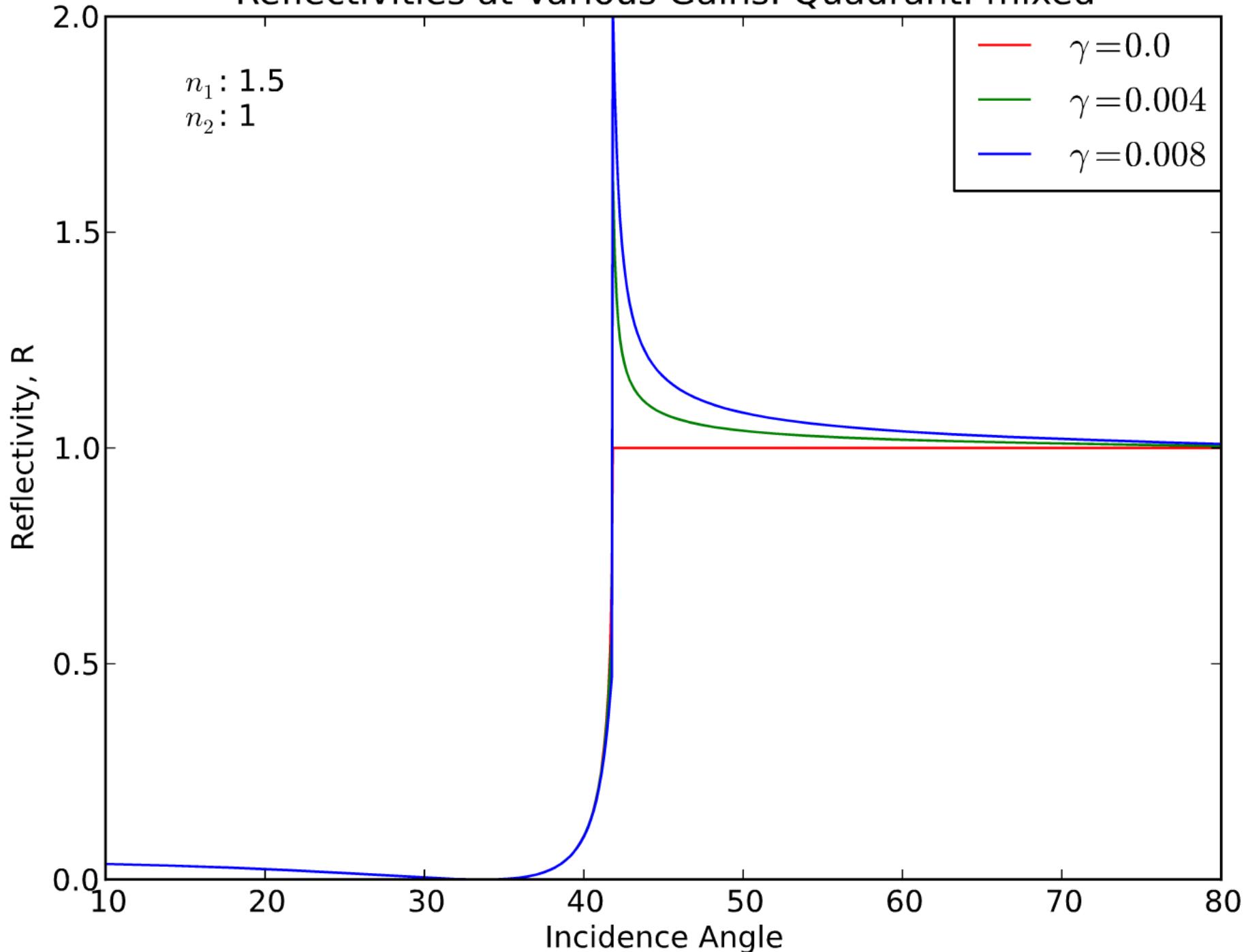
$$e^{i\tilde{k}_{tz}} = \begin{cases} e^{-i|k'_{tz}|z - |k''_{tz}|z} & 2^{nd} \text{Quadrant} \\ e^{i|k'_{tz}|z + |k''_{tz}|z} & 4^{th} \text{Quadrant} \end{cases}$$

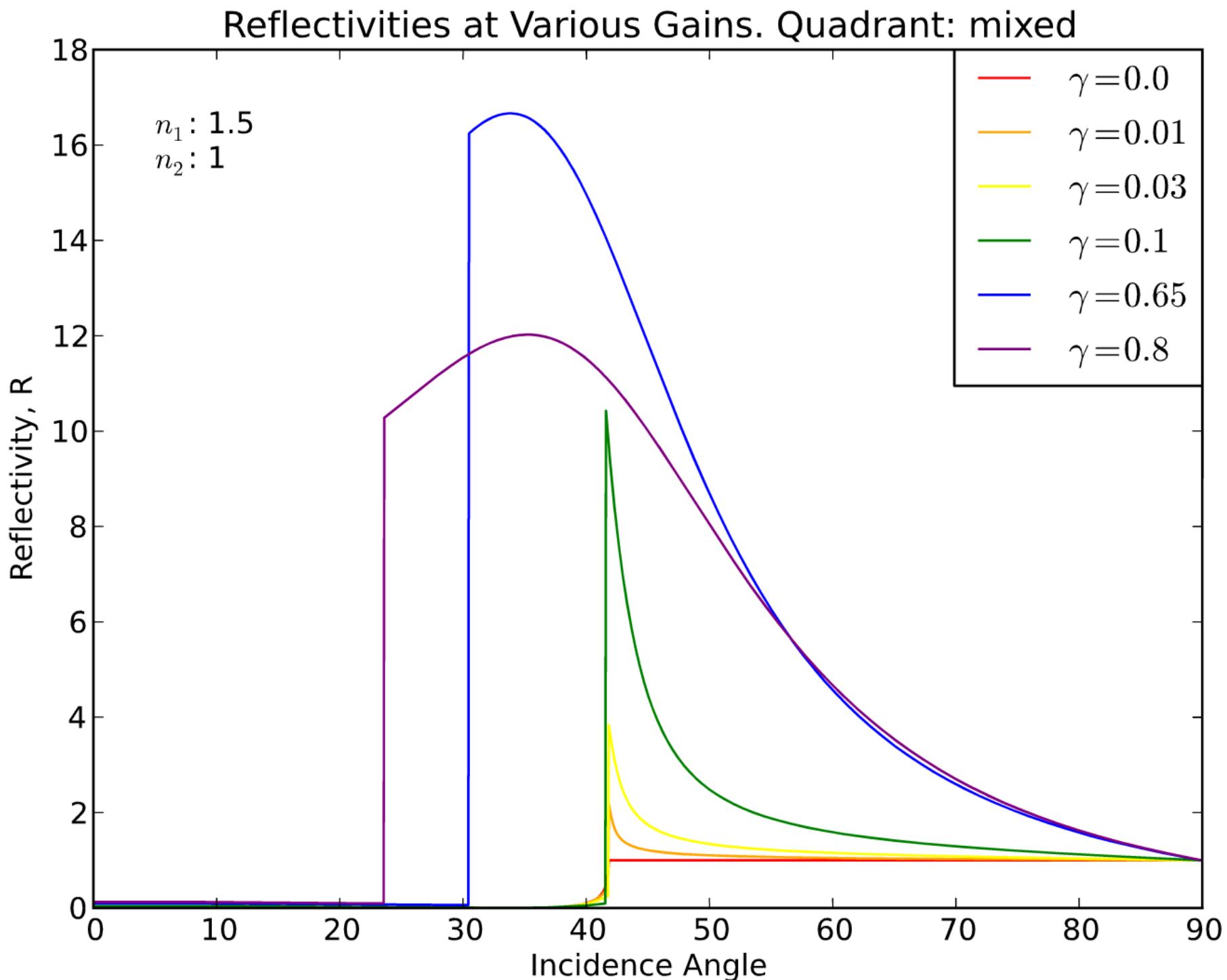
Gain Medium, Beyond Critical angle



$$e^{i\tilde{k}_{tz}} = \begin{cases} e^{-i|k'_{tz}|z - |k''_{tz}|z} & 2^{nd} \text{Quadrant} \\ e^{i|k'_{tz}|z + |k''_{tz}|z} & 4^{th} \text{Quadrant} \end{cases}$$

Reflectivities at Various Gains. Quadrant: mixed





Im

Lossy Medium
Beyond Critical Angle

Lossy Medium
Below Critical Angle

A'

B'

Re

D'

C'

Gainy Medium
Beyond Critical Angle

Gainy Medium
Below Critical Angle

D

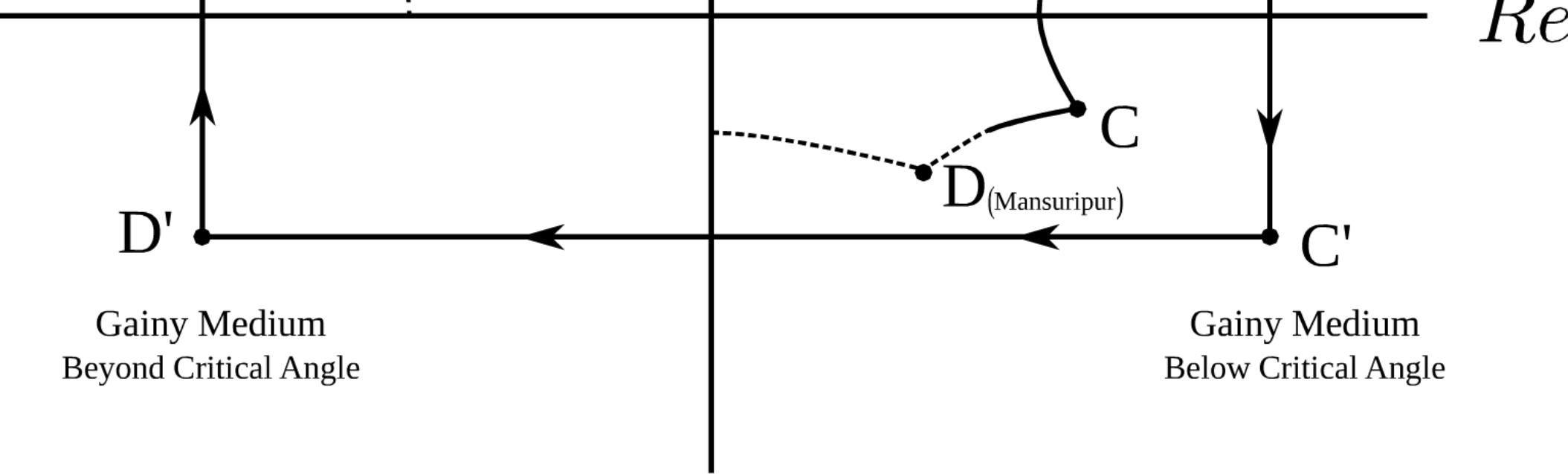
A

B

C

D_(Mansuripur)

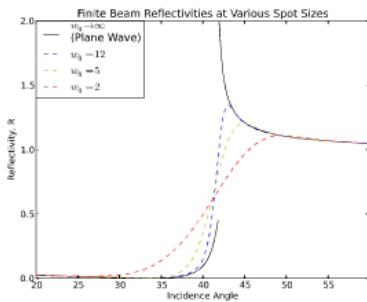
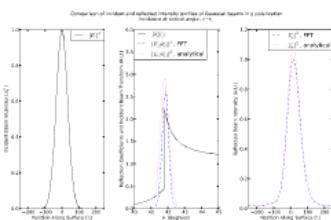
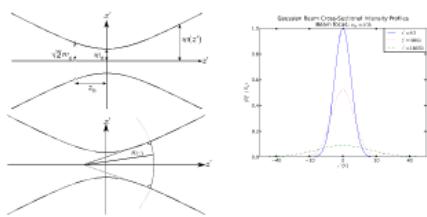
C
(Callary &
Carniglia)



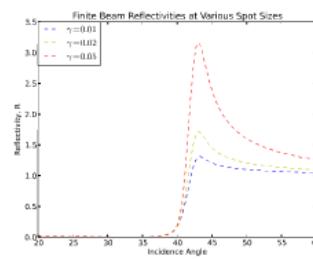
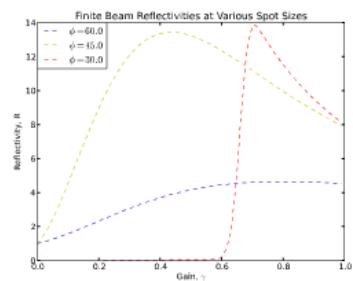
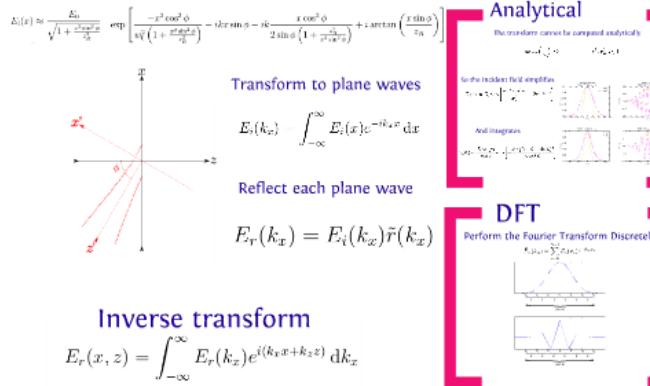
Finite Beam

A Gaussian Beam

$$E(x') = E_0 \frac{w_0}{w^2(z')} \exp \left[-\frac{x'^2 + y'^2}{w^2(z')} + ik \left(z' - \frac{x'^2 + y'^2}{2R(z')} \right) + i\xi(z') \right]$$

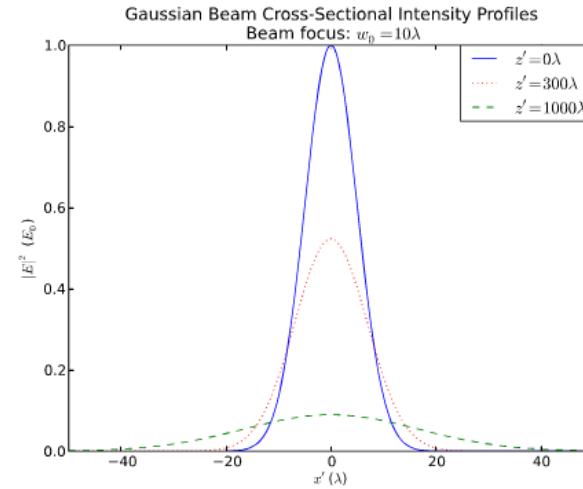
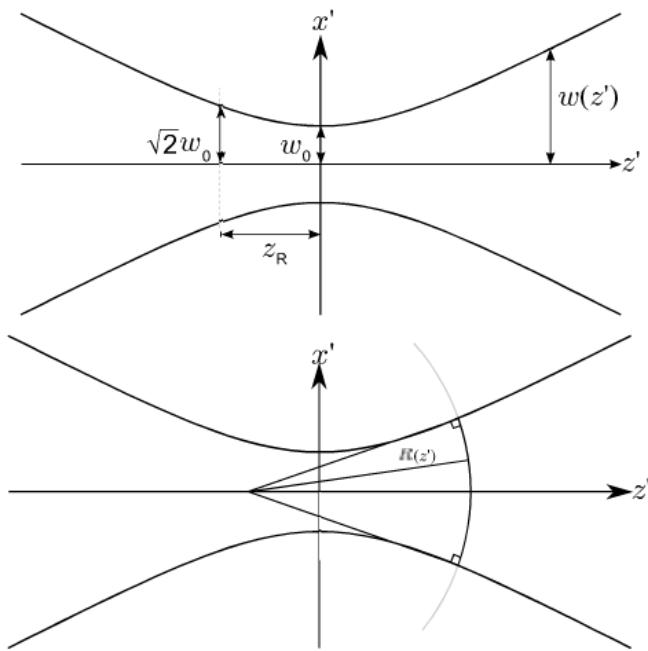


The Fourier Transform Method



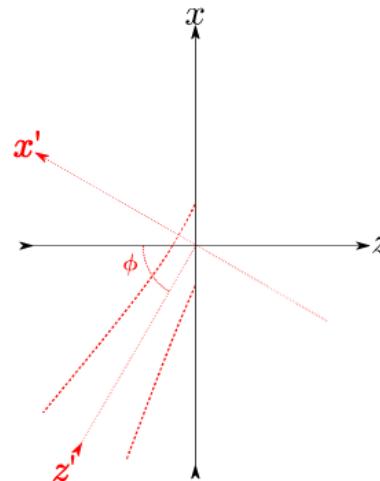
A Gaussian Beam

$$E(\mathbf{x}') = E_0 \frac{w_0}{w^2(z')} \exp \left[-\frac{x'^2 + y'^2}{w^2(z')} + ik \left(z' - \frac{x'^2 + y'^2}{2R(z')} \right) + i\xi(z') \right]$$



The Fourier Transform Method

$$E_i(x) \approx \frac{E_0}{\sqrt{1 + \frac{x^2 \sin^2 \phi}{z_R^2}}} \exp \left[\frac{-x^2 \cos^2 \phi}{w_0^2 \left(1 + \frac{x^2 \sin^2 \phi}{z_R^2} \right)} + ikx \sin \phi - ik \frac{x \cos^2 \phi}{2 \sin \phi \left(1 + \frac{z_R^2}{x^2 \sin^2 \phi} \right)} + i \arctan \left(\frac{x \sin \phi}{z_R} \right) \right]$$



Transform to plane waves

$$E_i(k_x) = \int_{-\infty}^{\infty} E_i(x) e^{-ik_x x} dx$$

Reflect each plane wave

$$E_r(k_x) = E_i(k_x) \tilde{r}(k_x)$$

Inverse transform

$$E_r(x, z) = \int_{-\infty}^{\infty} E_r(k_x) e^{i(k_x x + k_z z)} dk_x$$

Analytical

The transform cannot be computed analytically

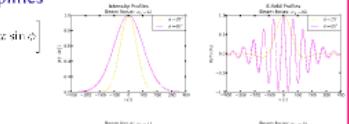
$$\arctan \left(\frac{z'}{x_0} \right) \approx 0 \quad x^2/z_R^2 \ll 1$$

So the incident field simplifies

$$E_i(x) \approx E_0 \exp \left[-\frac{x^2 \cos^2 \phi}{w_0^2} + ikx \sin \phi \right]$$

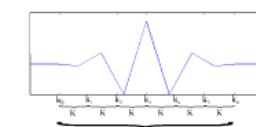
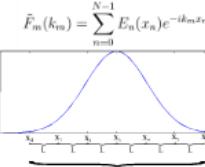
And integrates

$$E_r(k_x) \approx \frac{E_0 w_0 \sqrt{\pi}}{\cos \phi} \exp \left[-\frac{k^2 w_0^2 (\cos \phi - \sin \phi)^2}{4 \cos^2 \phi} \right]$$



DFT

Perform the Fourier Transform Discretely

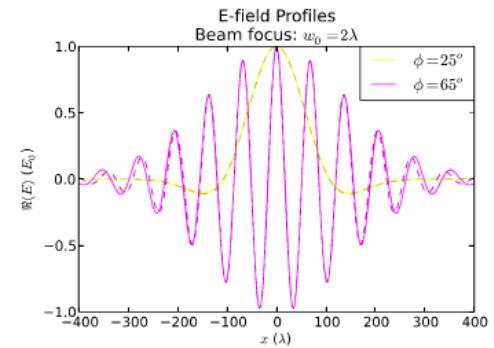
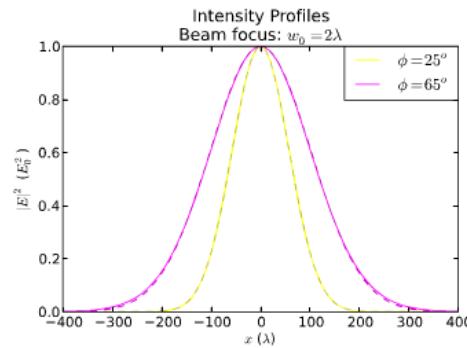


The transform cannot be computed analytically

$$\arctan\left(\frac{z'}{z_R}\right) \approx 0 \quad x^2/z_R^2 \ll 1$$

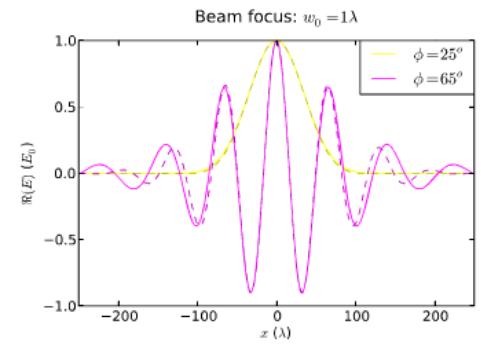
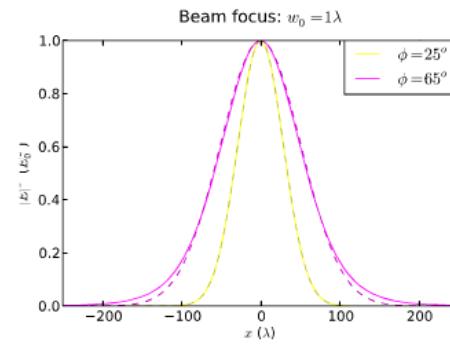
So the incident field simplifies

$$E_i(x) \approx E_0 \exp\left[\frac{-x^2 \cos^2 \phi}{w_0^2} + ikx \sin \phi\right]$$

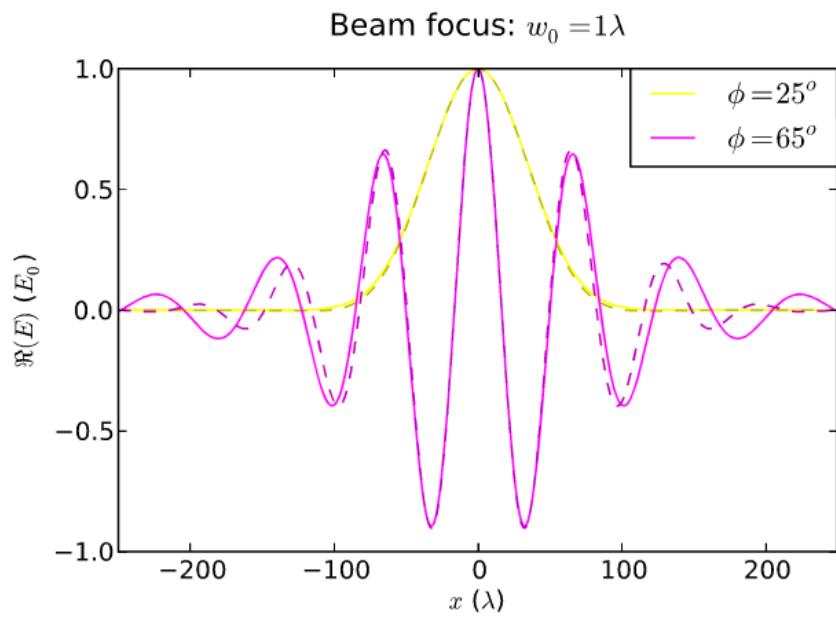
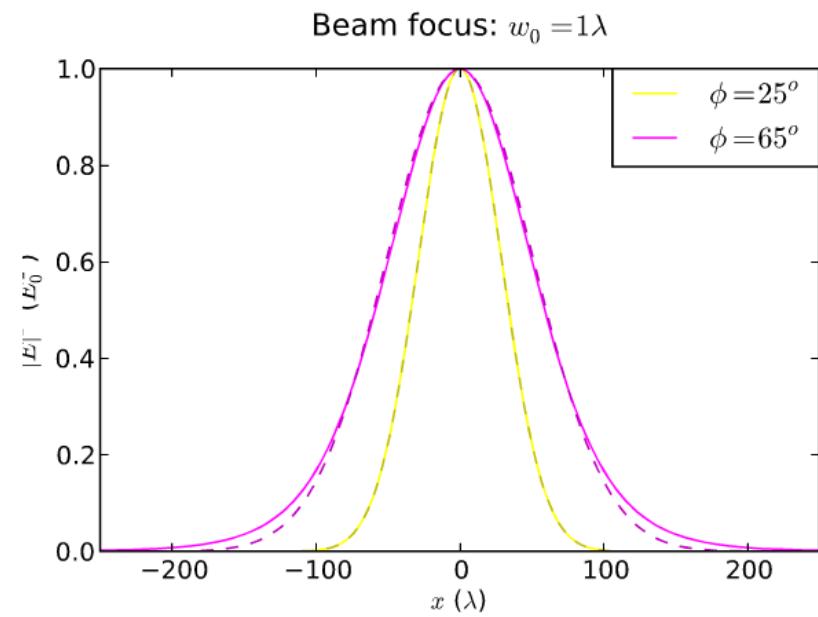
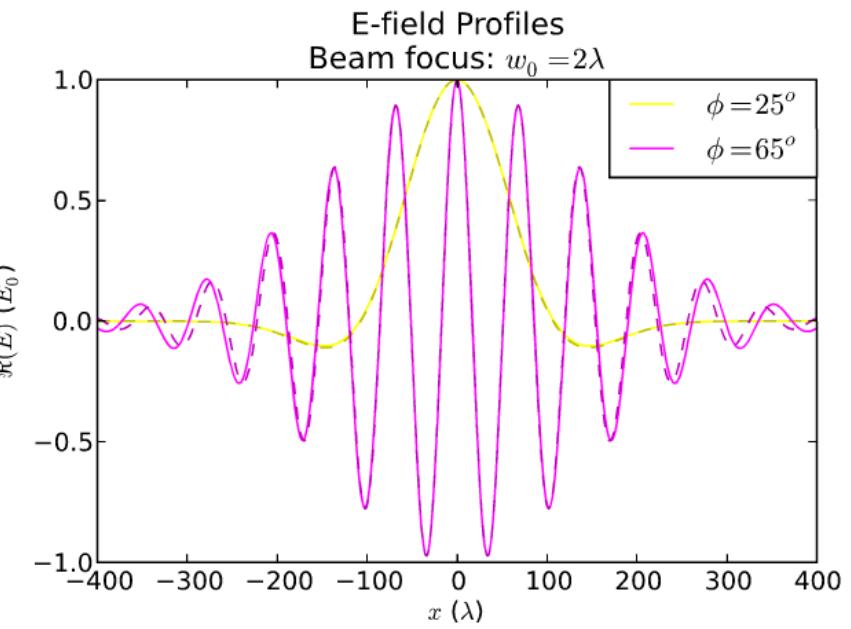
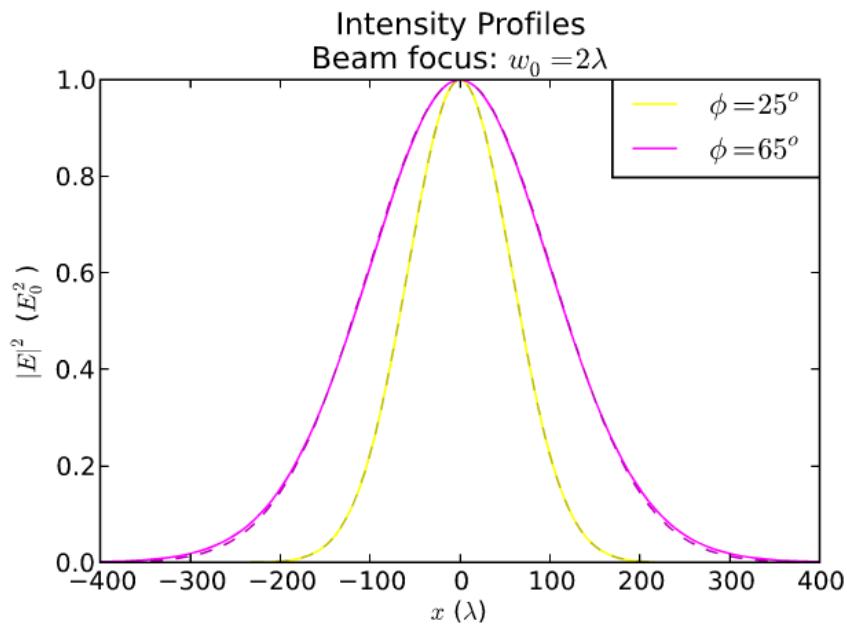


And integrates

$$E_i(\theta_i) \approx \frac{E_0 w_0 \sqrt{\pi}}{\cos \phi} \exp\left[-\frac{k^2 w_0^2 (\sin \phi - \sin \theta_i)^2}{4 \cos^2 \phi}\right]$$



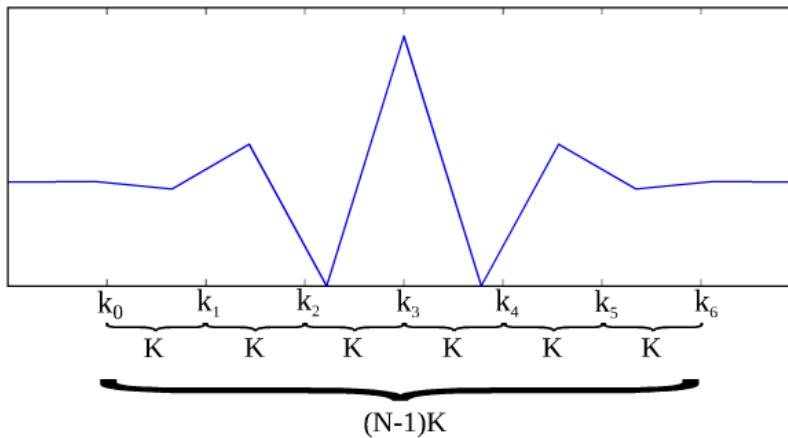
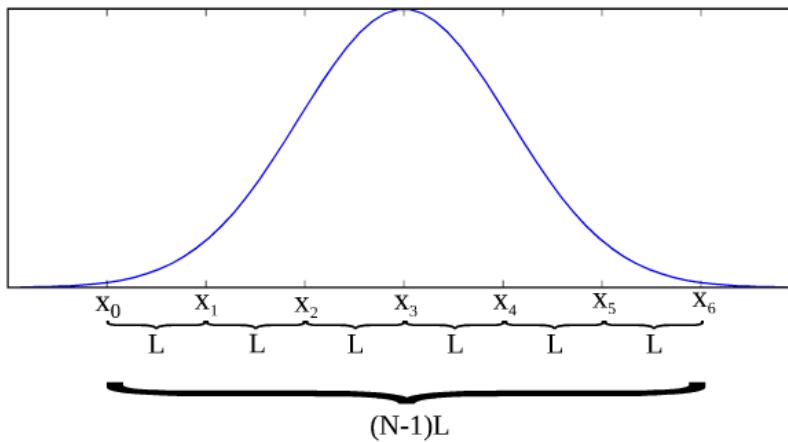
S



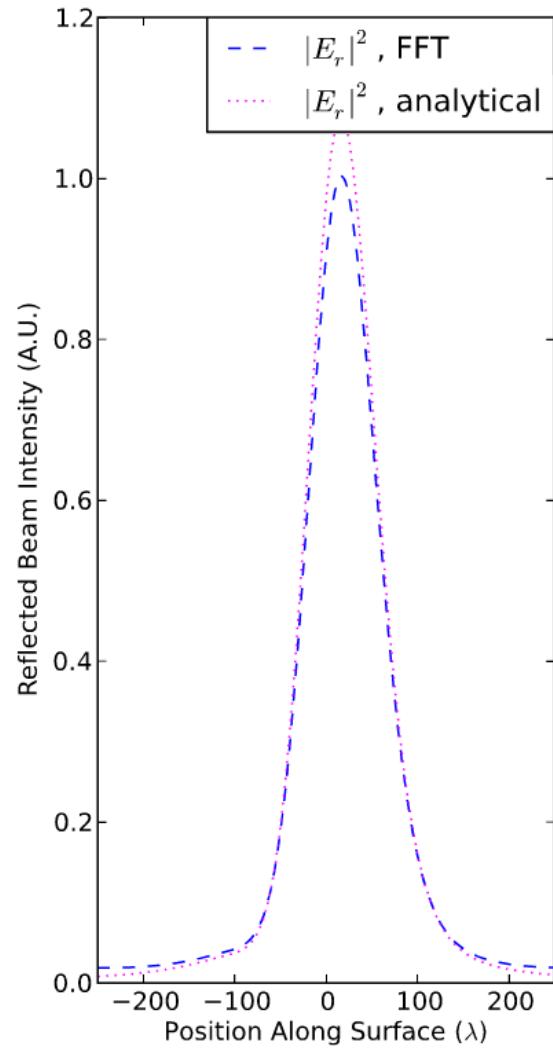
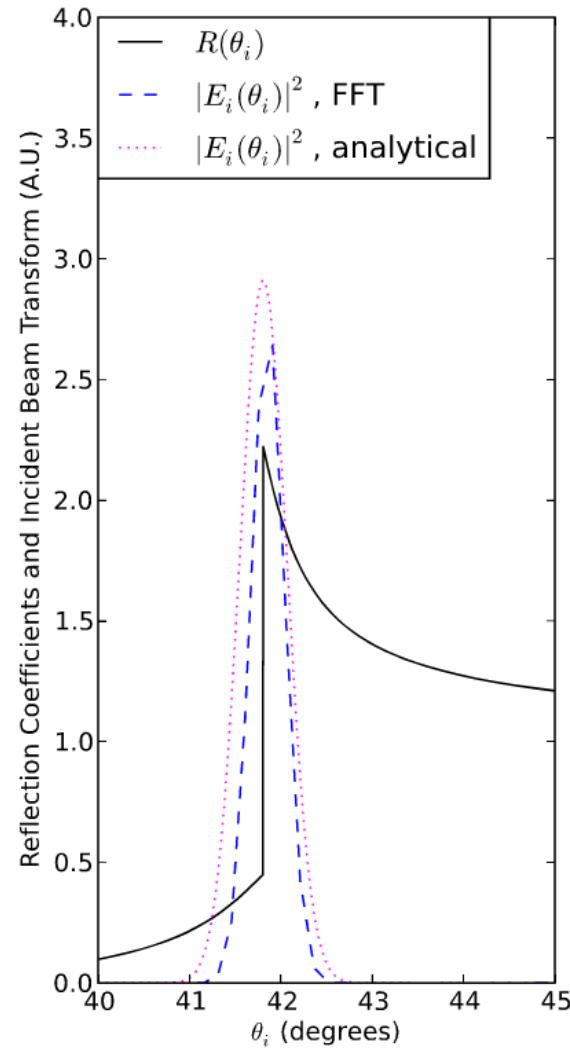
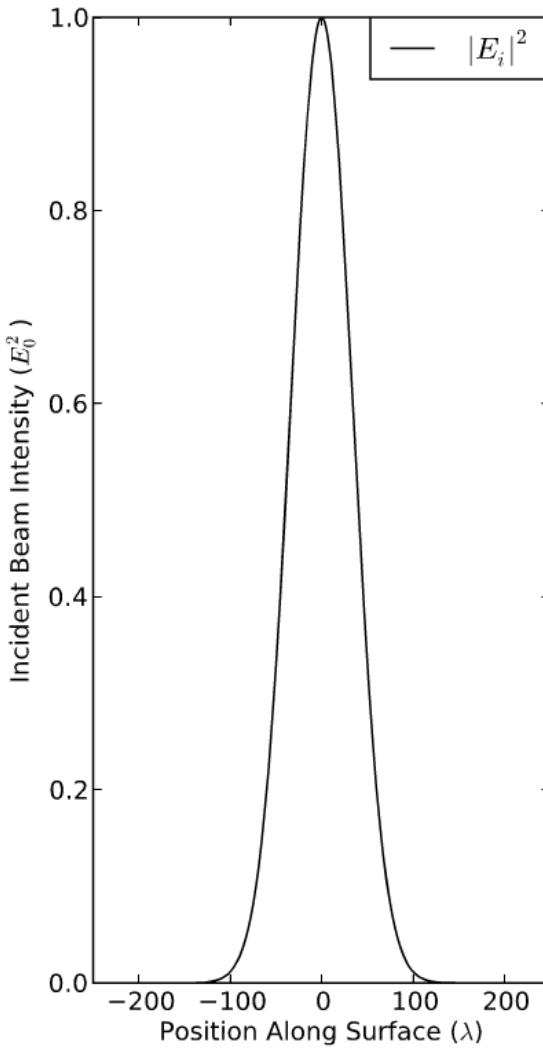
DFT

Perform the Fourier Transform Discretely

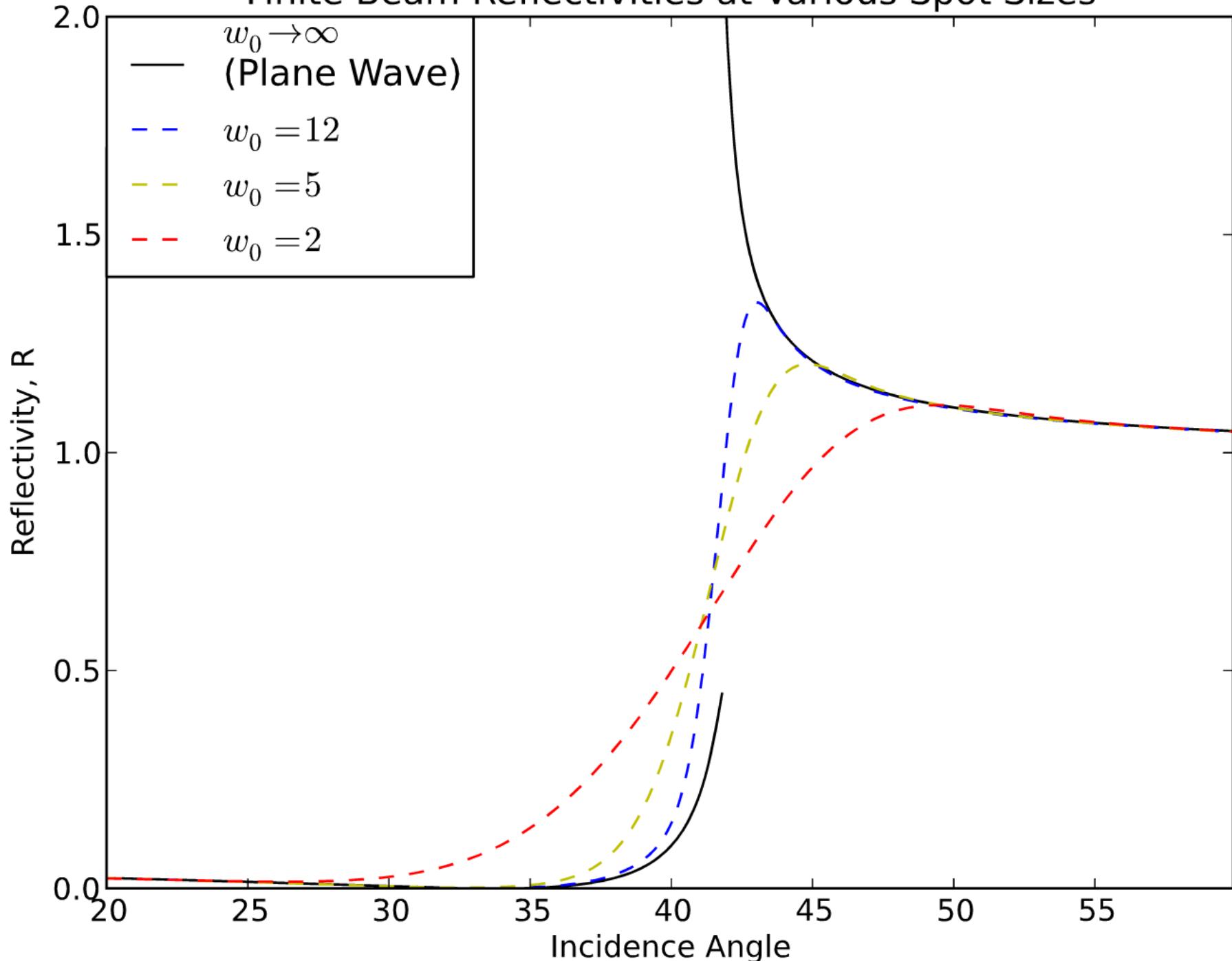
$$\tilde{F}_m(k_m) = \sum_{n=0}^{N-1} E_n(x_n) e^{-ik_m x_n}$$

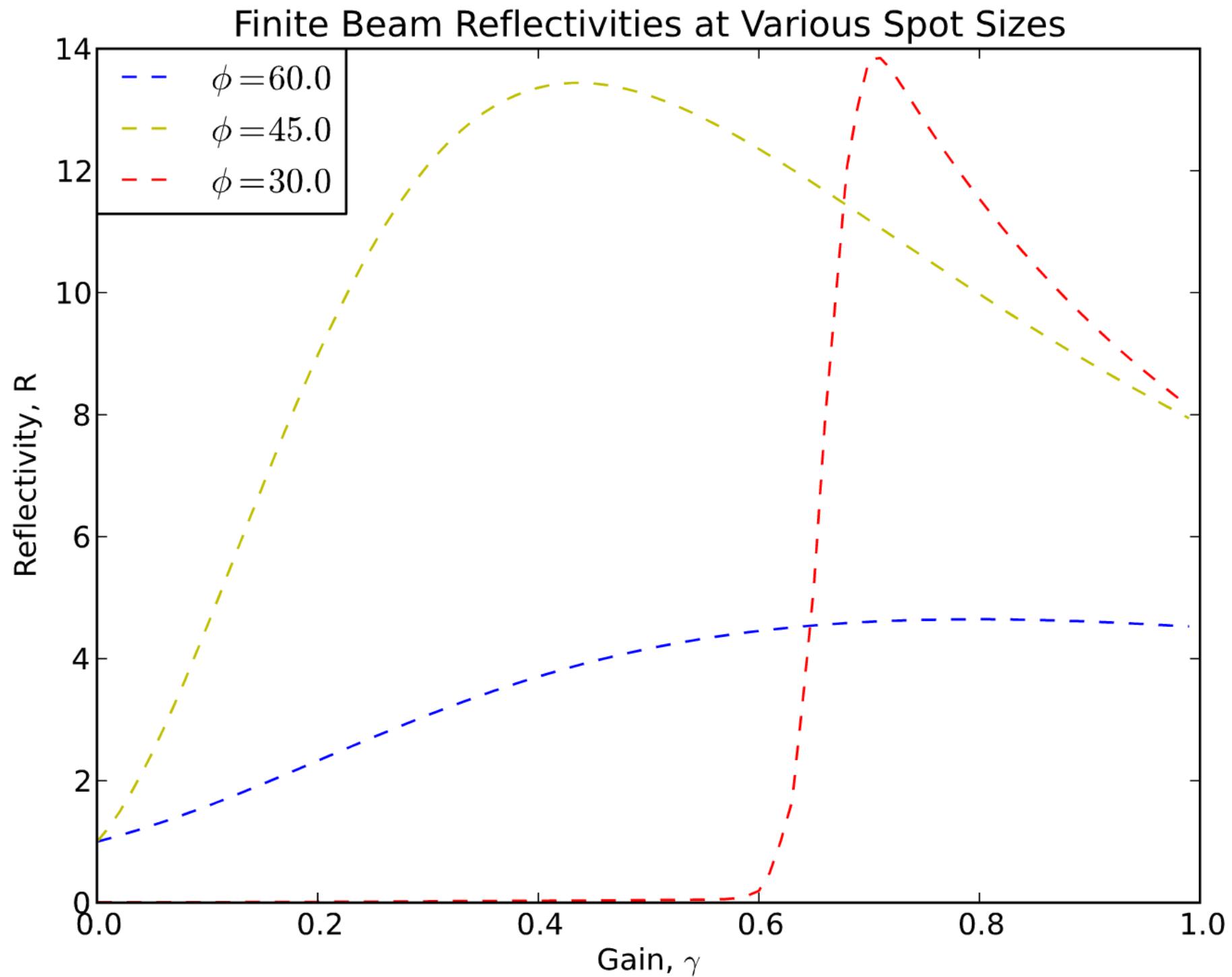


Comparison of incident and reflected intensity profiles of Gaussian beams in p polarization
 Incidence at critical angle, $\phi = \theta_c$

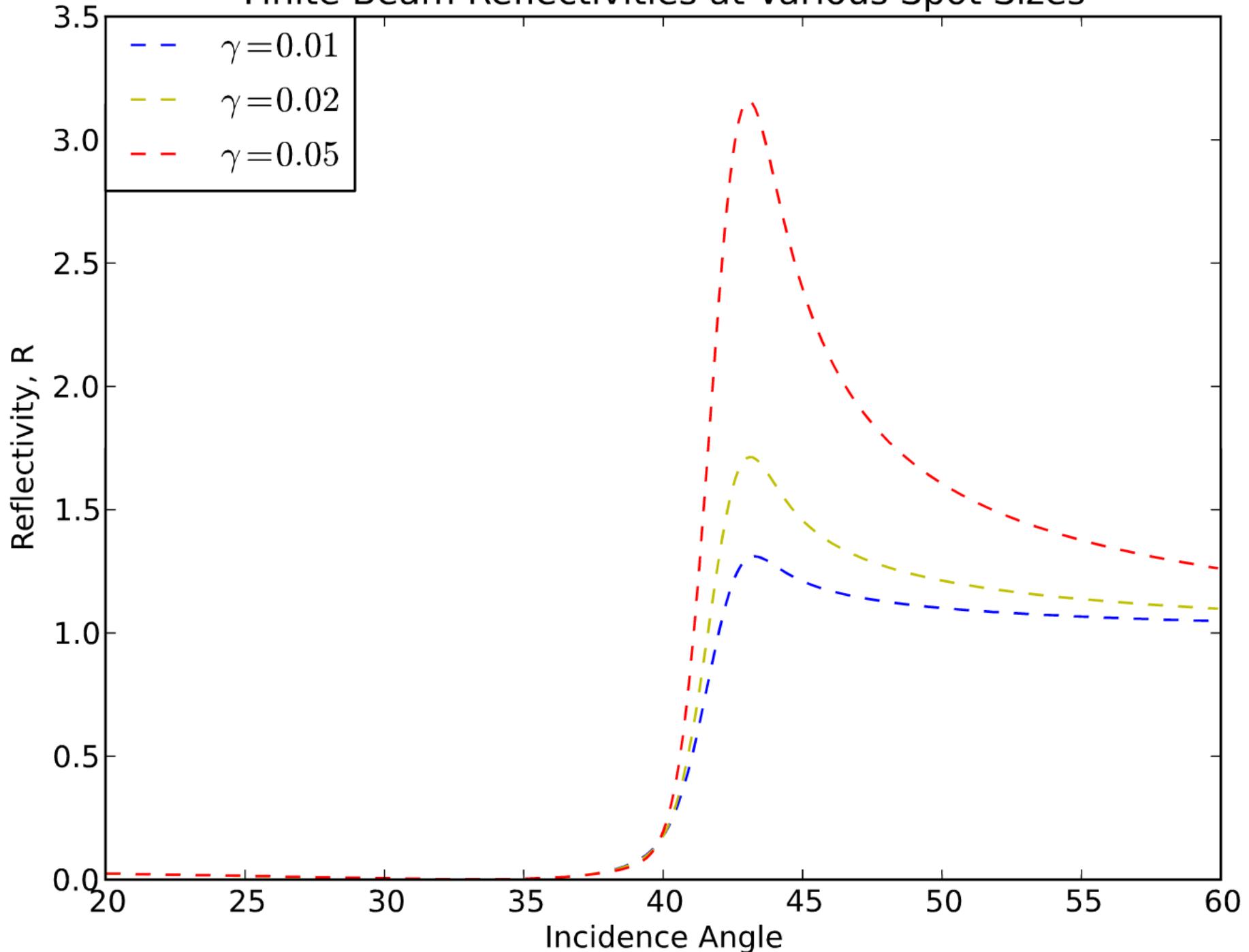


Finite Beam Reflectivities at Various Spot Sizes

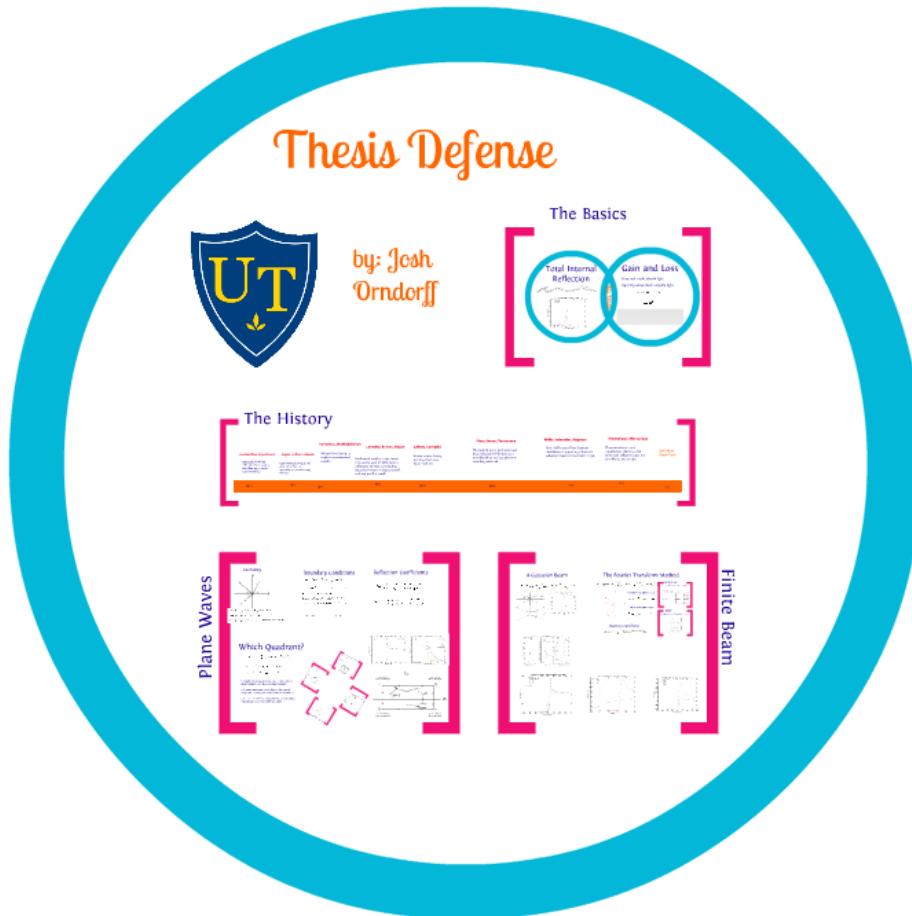




Finite Beam Reflectivities at Various Spot Sizes



Amplified Total Internal Reflection at the Surface of a Gain Medium



Special thanks to Dr. Deck,
Dr. Karpov, and Dr. Bagley